Quantum transport simulations using Kwant

Anton Akhmerov

ICTP Summer School "Advances in Condensed Matter Physics", 9 May 2019





I borrowed from a tutorial by J.B. Weston (CC-BY-SA), link at the end

- 1. How to solve a scattering problem?
- 2. How to define and solve it with Kwant?

Kwant

- ► An open source package for quantum transport
- https://kwant-project.org
- ▶ New version (1.4) released a couple of months ago
- Works with *tight binding* models
- Main focus on the scattering matrix formalism of quantum transport

Maintainers: Christoph Groth, Michael Wimmer, Anton Akhmerov, Xavier Waintal, Joseph Weston

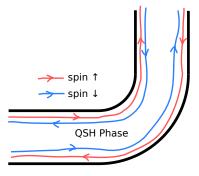
Contributors: Jörg Behrmann, Paul Clisson, Mathieu Istas, Daniel Jaschke, Bas Nijholt, Michał Nowak, Viacheslav Ostroukh, Pablo Pérez Piskunow, Tómas Örn Rosdahl, Sebastian Rubbert, Rafał Skolasiński, Adrien Sorgniard, Dániel Varjas, Thomas Kloss, Pierre Carmier

```
Supported by: ERC, NWO, and ONR
```

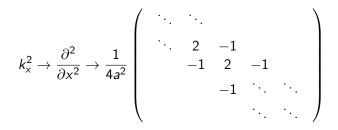
Case study: a bend in a quantum spin Hall insulator

- Topologically protected transmission across a bend
- $k \cdot p$ Hamiltonian

$$\begin{aligned} H(k_x, k_y) &= \\ &+ \mu - D(k_x^2 + k_y^2) \\ &+ [M - B(k_x^2 + k_y^2)]\tau_z \\ &+ A[k_x \sigma_z \tau_x + k_y \tau_y] \end{aligned}$$



- Kwant works with *tight-binding models* and a finite number of degrees of freedom
- \blacktriangleright \Rightarrow discretize space onto a square lattice with lattice spacing *a*



Discretize

```
import kwant
bhz_continuum = '''
+ mu * kron(sigma 0, sigma_0)
+ M * kron(sigma 0, sigma_z)
- B * (k_x**2 + k_y**2) * kron(sigma_0, sigma_z) - D * (k_x**2 + k_y**2) * kron(sigma_0, sigma_0)
+ A * (k_x * kron(sigma_z, sigma_x) + k_y * kron(sigma_0, sigma_y))
gshe model = kwant.continuum.discretize(bhz continuum)
```

Check the matrix elements

tb_matrix_elements, coords = kwant.continuum.discretize_symbolic(bhz_continuum)
tb_matrix_elements[0, 0] # onsite

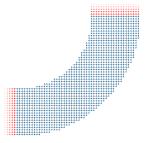
$\left[-\frac{2B}{a_{y}^{2}}-\frac{2B}{a_{x}^{2}}-\frac{2D}{a_{y}^{2}}-\frac{2D}{a_{x}^{2}}+M+\mu\right]$	0	0	0
0	$\frac{2B}{a_{y}^{2}} + \frac{2B}{a_{x}^{2}} - \frac{2D}{a_{y}^{2}} - \frac{2D}{a_{x}^{2}} - M + \mu$	0	0
0	0	$-\frac{2B}{a_{y}^{2}} - \frac{2B}{a_{x}^{2}} - \frac{2D}{a_{y}^{2}} - \frac{2D}{a_{x}^{2}} + M + \mu$	0
0	0	0	$\frac{2B}{a_{y}^{2}} + \frac{2B}{a_{x}^{2}} - \frac{2D}{a_{y}^{2}} - \frac{2D}{a_{x}^{2}} - M + \mu$

Check the symmetries

symmetries = kwant.qsymm.find_builder_symmetries(qshe_model)
print(f"{len(symmetries)} symmetries found.")

17 symmetries found.

Infinite system



has the Hamiltonian:

$$H = \begin{pmatrix} \ddots & V_L & \\ V_L^{\dagger} & H_L & V_L \\ & V_L^{\dagger} & H_L & V_L \\ & & V_L^{\dagger} & H_S \end{pmatrix}$$

 $(H_L \text{ and } V_L \text{ are block-diagonal if there are many leads})$

Scattering region

kwant.plot(syst);



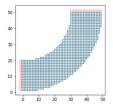
Make a lead

```
lead x = kwant.Builder(
    symmetry=kwant.TranslationalSymmetry((-1, 0)),
    time_reversal=np.kron(1j * sigma_y, sigma_0),
    conservation_law=np.kron(-sigma_z, sigma_0), # spin conservation
)
lead_x.fill(qshe_model, lambda site: 1 <= site.tag[1] <= 20, (0, 1))</pre>
```

Combine everything

```
syst.attach_lead(lead_x)
syst.attach_lead(lead_y)
```

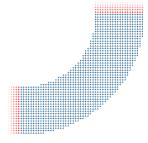
kwant.plot(syst);



fsyst = syst.finalized() # Transform the system into an efficient form for numerics
print(fsyst)

<FiniteSystem with 1272 sites, 4888 hoppings, and parameters: ('M', 'B', 'D', 'A', 'mu')>

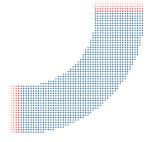
Scattering problem



Now to get conductance and other observables, we only need to calculate Σ_{lead} , G^R , $G^<$, and we're done.

(that is what one mostly hears)

Scattering problem

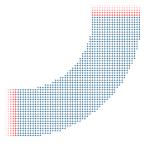


Now to get conductance and other observables, we only need to calculate Σ_{lead} , G^R , $G^<$, and we're done.

(that is what one mostly hears)

... But there is a more intuitive definition.

Scattering problem



Problem statement:

$$H = \begin{pmatrix} \ddots & V_L & & \\ V_L^{\dagger} & H_L & V_L & \\ & V_L^{\dagger} & H_L & V_L \\ & & V_L^{\dagger} & H_S \end{pmatrix}, \quad \psi = \begin{pmatrix} \vdots \\ \psi_2 \\ \psi_1 \\ \psi_1 \\ \psi_S \end{pmatrix}$$

We are solving $(H - E)\psi = 0$ with a fixed E.

Since the leads are translationally invariant, we decompose the lead wave function into eigenvectors of translation (lead modes)

$$\psi_j = \lambda^j \psi_0,$$

Since the leads are translationally invariant, we decompose the lead wave function into eigenvectors of translation (lead modes)

$$\psi_j = \lambda^j \psi_0,$$

Now substitute into *H*:

$$V_L\psi_0 + H\lambda\psi_0 + V_L^{\dagger}\lambda^2\psi_0 = 0,$$

Since the leads are translationally invariant, we decompose the lead wave function into eigenvectors of translation (lead modes)

$$\psi_j = \lambda^j \psi_0,$$

Now substitute into *H*:

$$V_L\psi_0 + H\lambda\psi_0 + V_L^{\dagger}\lambda^2\psi_0 = 0,$$

or in a matrix form:

$$\begin{pmatrix} -V_L^{-1}H_L & -V_L^{-1} \\ V_L^{\dagger} & 0 \end{pmatrix} \begin{pmatrix} \psi_0 \\ V_L^{\dagger}\psi_1 \end{pmatrix} = \lambda^{-1} \begin{pmatrix} \psi_0 \\ V_L^{\dagger}\psi_1 \end{pmatrix}$$

Since the leads are translationally invariant, we decompose the lead wave function into eigenvectors of translation (lead modes)

$$\psi_j = \lambda^j \psi_0,$$

Now substitute into *H*:

$$V_L\psi_0 + H\lambda\psi_0 + V_L^{\dagger}\lambda^2\psi_0 = 0,$$

or in a matrix form:

$$\begin{pmatrix} -V_L^{-1}H_L & -V_L^{-1} \\ V_L^{\dagger} & 0 \end{pmatrix} \begin{pmatrix} \psi_0 \\ V_L^{\dagger}\psi_1 \end{pmatrix} = \lambda^{-1} \begin{pmatrix} \psi_0 \\ V_L^{\dagger}\psi_1 \end{pmatrix}$$

... or in a more stable form and reduced basis:

$$\begin{pmatrix} iA^{\dagger}\tilde{H}^{-1}B & -A^{\dagger}\tilde{H}^{-1}B \\ -1 + iB^{\dagger}\tilde{H}^{-1}B & -B^{\dagger}\tilde{H}^{-1}B \end{pmatrix} \begin{pmatrix} \tilde{\psi}_{0} \\ \tilde{\psi}_{1} \end{pmatrix} = \lambda^{-1} \begin{pmatrix} A^{\dagger}\tilde{H}^{-1}A & -1 - iA^{\dagger}\tilde{H}^{-1}A \\ B^{\dagger}\tilde{H}^{-1}A & -iB^{\dagger}\tilde{H}^{-1}A \end{pmatrix} \begin{pmatrix} \tilde{\psi}_{0} \\ \tilde{\psi}_{1} \end{pmatrix}$$
$$\tilde{H} = H_{L} + iAA^{\dagger} + iBB^{\dagger}, \quad V_{L} = AB^{\dagger}, \quad \tilde{\psi}_{0} = B^{+}\psi_{0}, \quad \tilde{\psi}_{1} = A^{+}\psi_{1}$$

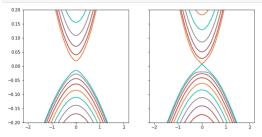
Split all the eigenvectors into incoming, outgoing and evanecsent, such that:

$$egin{aligned} &\langle\psi_{
m in}|\hat{I}|\psi_{
m in}
angle = 1 \ &\langle\psi_{
m out}|\hat{I}|\psi_{
m out}
angle = -1 \ &\langle\psi_{
m evan}|\hat{I}|\psi_{
m out}
angle = 0, \quad |\lambda_{
m evan}| < 1 \end{aligned}$$

Dispersion

topo_params = dict(A=0.09, B=-0.18, D=-0.065, M=-0.02, mu=0) # Parameters for topological phase trivial_params = dict(A=0.09, B=-0.18, D=-0.065, M=0.01, mu=0) # Parameters for trivial phase

fig, (ax1, ax2) = plt.subplots(1, 2, figsize=(10, 5), sharey=True)
kwant.plotter.bands(flead x, momenta=np.linspace(-2, 2, 101), params=trivial_params, show=False, ax=ax1)
kwant.plotter.bands(flead_x, momenta=np.linspace(-2, 2, 101), params=topo_params, show=False, ax=ax2)
ax1.set_ylim(-.2, .2);



Modes at E = 0

propagating_modes, _ = flead_x.modes(energy=0, params=topo_params)

```
print(propagating_modes.wave_functions.shape)
print(propagating_modes.velocities)
print(propagating_modes.momenta)
```

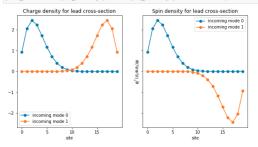
```
(80, 4)
[-0.08354984 -0.08354984 0.08354984 0.08354984]
[ 0.08609145 0.08609145 -0.08609145 -0.08609145]
```

Mode wave functions

```
phi0, phi1 = propagating_modes.wave_functions[:, :2].transpose()
```

density = kwant.operator.Density(fsyst.leads[0]) # calculate $|\varphi_i|^{*2}$ for each site (summing degrees of freedom) spin_density = kwant.operator.Density(fsyst.leads[0], np.kron(sigma_z, sigma_0))

fig, (ax1, ax2) = plt.subplots(1, 2, figsize=(10, 5), sharey=True)
plot_densities(density(phi0), density(phi1), title='Charge', ax=ax1)
plot_densities(spin_density(phi10), spin_density(phi11), title='Spin', ax=ax2)



Substitute the lead modes into the Hamiltonian and get a linear system:

$$\begin{pmatrix} -U_{\rm out} & 1\\ V_L^{\dagger} U_{\rm out} \Lambda_{\rm out} & H_{\rm S} \end{pmatrix} \begin{pmatrix} {\bf S}\\ \psi_{\mathcal{S}} \end{pmatrix} = \begin{pmatrix} U_{\rm in}\\ -V_{\rm L}^{\dagger} U_{\rm in} \Lambda_{\rm in}, \end{pmatrix}$$
(1)

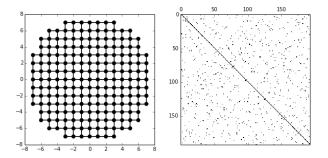
with $U_{\rm in}$ and $U_{\rm out}$ wave functions of incoming and outgoing modes, and $\Lambda \equiv {\rm diag}(\lambda_i)$.

Substitute the lead modes into the Hamiltonian and get a linear system:

$$\begin{pmatrix} -U_{\rm out} & 1\\ V_L^{\dagger} U_{\rm out} \Lambda_{\rm out} & H_{\rm S} \end{pmatrix} \begin{pmatrix} {\bf S}\\ \psi_{\cal S} \end{pmatrix} = \begin{pmatrix} U_{\rm in}\\ -V_{\rm L}^{\dagger} U_{\rm in} \Lambda_{\rm in}, \end{pmatrix}$$
(1)

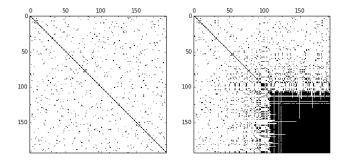
with $U_{\rm in}$ and $U_{\rm out}$ wave functions of incoming and outgoing modes, and $\Lambda \equiv {\rm diag}(\lambda_i)$.

Next: write down these linear equations and solve them. (NB: if we start by eliminating **S**, the rhs becomes $H_S + \Sigma$) The scattering region is large and its Hamiltonian is sparse.

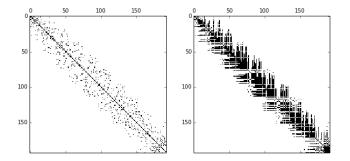


Solving a scattering problem

Naive solution: gaussian elimination (dense LU decomposition, cost L^6)

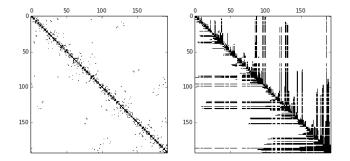


Better option: reshuffle to reduce bandwidth ("recursive Green's functions", cost L^4)



Solving a scattering problem

What specialized libraries do (nested dissection, cost L^3 , $\sim 10 \times$ better):



Scattering matrix

S = kwant.smatrix(fsyst, energy=0, params=topo_params)
print(np.round(S.data, 2))
print(S.transmission(1, 0)) # Tr(t . t^t)
[[-0. -0.j 0. +0.j 0.27-0.96j 0. +0.j]
[-0. +0.j 0. -0.j -0. +0.j 0.5 +0.86j]
[0.5 +0.66j 0. +0.j -0. -0.j 0. +0.j 1]
[-0. +0.j 0.27-0.96j -0. +0.j -0. +0.j]
[.999995792536172

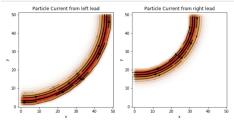
Scattering matrix

```
S = kwant.smatrix(fsyst, energy=0, params=topo_params)
print(np.round(S.data, 2))
print(S.transmission(1, 0)) # Tr(t . t^t)
[[-0. -0.j 0. +0.j 0.27-0.96j 0. +0.j ]
[-0. +0.j 0. -0.j -0. +0.j 0.5 +0.86j]
[0.5 +0.66j 0. +0.j -0. -0.j 0. +0.j 1]
[-0. +0.j 0.27-0.96j -0. +0.j -0. +0.j ]
[.999995792536172
```

Current density

```
current = kwant.operator.Current(fsyst).bind(params=topo_params)
```

```
(ax0, ax1) = prepare_axes('Particle Current');
kwant.plotter.current(fsyst, current(sl0) , ax=ax0)
kwant.plotter.current(fsyst, current(sr0), ax=ax1)
```



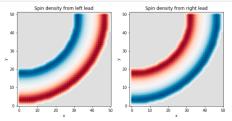
Scattering matrix

```
S = kwant.smatrix(fsyst, energy=0, params=topo_params)
print(np.round(S.data, 2))
print(S.transmission(1, 0)) # Tr(t . t^†)
[[-0. -0.j 0. +0.j 0.27-0.96j 0. +0.j ]
[-0. +0.j 0. -0.j 0. +0.j 0.5 +0.66j]
[0.5 +0.66j 0. +0.j -0. -0.j 0. +0.j 0. +0.j ]
[-0. +0.j 0.27-0.96j -0. +0.j 0. +0.j 1]
[.99999579253612
```

Spin density

```
spin density = kwant.operator.Density(fsyst, spinz operator).bind(params=topo params)
```

```
(ax0, ax1) = prepare_axes('Spin_density')
kwant.plotter.density(fsyst, spin_density(sl0) + spin_density(sl1), cmap='RdBu_r', ax=ax0)
kwant.plotter.density(fsyst, spin_density(sr0) + spin_density(sr1), cmap='RdBu_r', ax=ax1)
```



- Write down a problem + feed it to a good solver = problem solved
- Kwant does this with quantum transport
- Try a live version at https://tiny.cc/maryland-kwant-tutorial (I used code and images from it)

The end. Questions?